The *p*-adic Mehta Integral: Formulas, Functional Equations, and Combinatorics

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② Fix a measurable function E : ℝ^N → [-∞, ∞] that assigns each microstate a total energy E(x).

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- **3** Fix the Boltzmann constant k > 0 that makes $\frac{E(\mathbf{x})}{kT}$ dimensionless and define the **inverse temperature parameter** $\beta = \frac{1}{kT}$.

The energy E induces a probability distribution on the microstates:

$$d\mathbb{P}_{eta}(oldsymbol{x}) = rac{1}{\mathcal{Z}_{N}(eta)}e^{-eta E(oldsymbol{x})}doldsymbol{x} \quad ext{where} \quad \mathcal{Z}_{N}(eta) = \int_{\mathbb{R}^{N}}e^{-eta E(oldsymbol{x})}\,doldsymbol{x}$$

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• Intuition: Low-energy states are more probable than high-energy states. This disparity becomes more pronounced as $T \searrow 0$.

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- Important task: Determine the domain and explicit form of the canonical partition function Z_N .

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$$\mathcal{Z}_{\mathcal{N}}(eta) = \int_{\mathbb{R}^{\mathcal{N}}} e^{-rac{1}{2} \|oldsymbol{x}\|^2} \prod_{i < j} |x_i - x_j|^eta \, doldsymbol{x}$$

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- harmonic potential energies $\frac{1}{2\beta}x_i^2$ for i = 1, 2, ..., N and
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Conjecture (Mehta and Dyson, early 1960's)

$$\mathcal{Z}_{N}(\beta) = (2\pi)^{N/2} \prod_{j=1}^{N} \frac{\Gamma(1+j\beta/2)}{\Gamma(1+\beta/2)} \quad if \quad \operatorname{Re}(\beta) > -2/N$$

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Theorem (Bombieri, late 1970's)

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p-adic log-Coulomb gas

• Suppose the charges have random locations $x_1, \ldots, x_N \in \mathbb{Q}_p$ instead.

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- Now Q^N_p is the space of microstates *x* = (x₁,..., x_N) with standard norm || · ||_p and Haar measure *dx* defined by

$$\|\mathbf{x}\|_{p} = \max_{1 \le i \le N} |x_{i}|_{p}$$
 and $\int_{\mathbb{Z}_{p}^{N}} d\mathbf{x} = 1$

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• Choose an analogue $V(\mathbf{x})$ of the total harmonic potential, so that $e^{-\beta V(\mathbf{x})} = \rho(\|\mathbf{x}\|_p)$ is "nice" (like $e^{-\frac{1}{2}\|\mathbf{x}\|^2}$ for $\mathbf{x} \in \mathbb{R}^N$) and define

$$E(\mathbf{x}) = V(\mathbf{x}) - \sum_{i < j} \log |x_i - x_j|_p$$

Main question:

$$\mathcal{Z}_{N}(\beta) = \int_{\mathbb{Q}_{p}^{N}} \rho(\|\boldsymbol{x}\|_{p}) \prod_{i < j} |x_{i} - x_{j}|_{p}^{\beta} d\boldsymbol{x} = ???$$

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• Nice fact 1: It suffices to compute $\int_{\mathbb{Z}_p^N} \prod_{i < j} |x_i - x_j|_p^\beta dx$ because

$$\mathcal{Z}_{N}(\beta) = \left(\sum_{m \in \mathbb{Z}} \rho(p^{m}) p^{m(N + \binom{N}{2}\beta)}\right) \cdot (1 - p^{-(N + \binom{N}{2}\beta)}) \cdot \int_{\mathbb{Z}_{p}^{N}} \prod_{i < j} |x_{i} - x_{j}|_{p}^{\beta} d\boldsymbol{x}$$

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Nice fact 2: V₀ := {x ∈ Z^N_p : x_i = x_j for some i < j} has measure 0, so we only need to do the integral over Z^N_p \ V₀.

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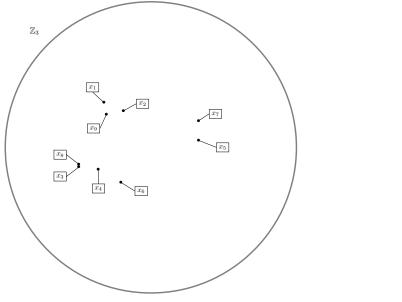
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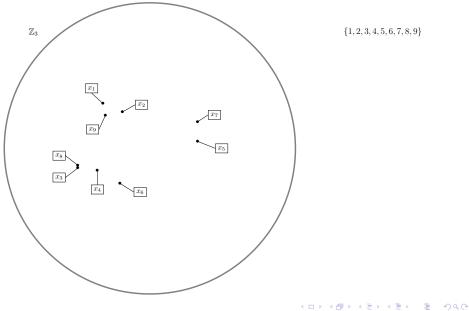
• Question: What do microstates $x \in \mathbb{Z}_p^N \setminus V_0$ look like?

What a microstate $\pmb{x} \in \mathbb{Z}_3^9 \setminus V_0$ looks like

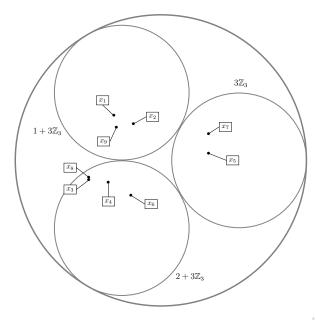


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What a microstate $x \in \mathbb{Z}_3^9 \setminus V_0$ looks like ...mod 3⁰



What a microstate $x \in \mathbb{Z}_3^9 \setminus V_0$ looks like ...mod 3^1

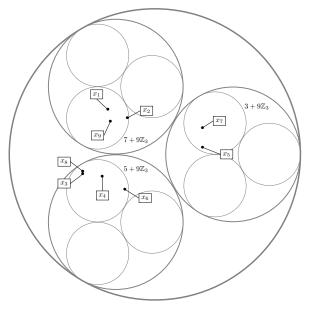


 $\{1, 2, 3, 4, 5, 6, 7, 8, 9\}$

 $\{5,7\}\ \{1,2,9\}\ \{3,4,6,8\}$

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What a microstate $x \in \mathbb{Z}_3^9 \setminus V_0$ looks like ...mod 3^2

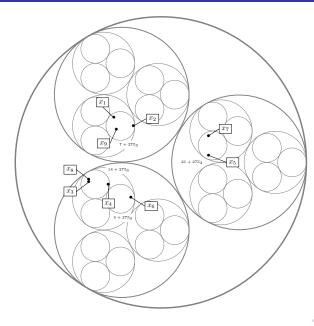


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What a microstate $x \in \mathbb{Z}_3^9 \setminus V_0$ looks like ...mod 3³



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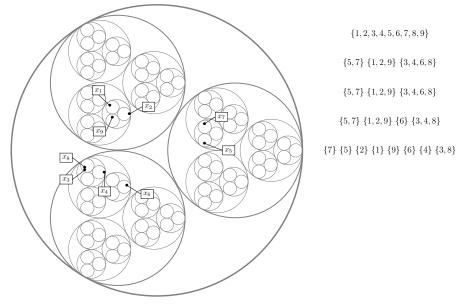
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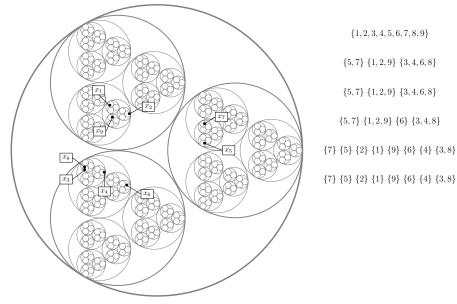
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What a microstate $x \in \mathbb{Z}_3^9 \setminus V_0$ looks like ...mod 3^4



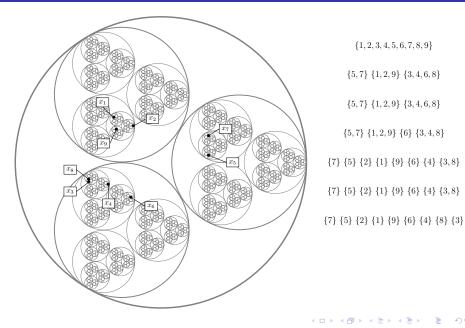
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What a microstate $x \in \mathbb{Z}_3^9 \setminus V_0$ looks like ...mod 3^5

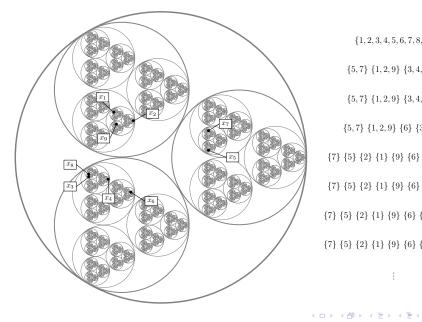


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What a microstate $x \in \mathbb{Z}_3^9 \setminus V_0$ looks like... mod 3^6







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 $\mathbf{h} = (h_0, h_1, h_2, h_3, h_4)$ and $\mathbf{n} = (n_0, n_1, n_2, n_3)$:

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$$h_0 = \{1, 2, 3, 4, 5, 6, 7, 8, 9\} \quad \text{appeared } n_0 = 1 \text{ time}$$

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What does the diagram tell us?

• The microstate $\pmb{x} \in \mathbb{Z}_3^9 \setminus V_0$ determines a pair of tuples

$$\mathbf{h} = (h_0, h_1, h_2, h_3, h_4) \quad \text{and} \quad \mathbf{n} = (n_0, n_1, n_2, n_3):$$

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Note: p = 3 is not special here. Many x in Z⁹₅, Z⁹₇, Z⁹₁₁, ..., etc., determine the same pair (𝑘, n) in the same way.

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- Note: p = 3 is not special here. Many x in Z⁹₅, Z⁹₇, Z⁹₁₁, ..., etc., determine the same pair (𝑘, n) in the same way.
- For any p, let $\mathcal{T}_p(\mathbf{fh}, \mathbf{n})$ be the set of all $\mathbf{x} \in \mathbb{Z}_p^9$ that determine $(\mathbf{fh}, \mathbf{n})$.

The value of $\prod_{i < j} |x_i - x_j|_{ ho}$ on $\mathcal{T}_{ ho}(\mathbf{fh}, oldsymbol{n})$

The value of $\prod_{i < j} |x_i - x_j|_p$ on $\mathcal{T}_p(\mathbf{fh}, \mathbf{n})$

• If $\mathcal{T}_{p}(\Uparrow, \textbf{\textit{n}}) \neq \varnothing$, then for every $\textbf{\textit{x}} \in \mathcal{T}_{p}(\Uparrow, \textbf{\textit{n}})$ we have

$$|x_i - x_j|_p = p^{1 - (n_0 + n_1 + \dots + n_{\ell_{ij}})}$$
 for $1 \le i < j \le 9$,

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where $\ell_{ij} = \max\{\ell : i \text{ and } j \text{ are in a common } \lambda \in \uparrow_{\ell}\}.$

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• This means any function that factors through $\mathbf{x} \mapsto (|x_i - x_j|_p)_{i < j}$ is constant on $\mathcal{T}_p(\mathbf{h}, \mathbf{n})$, with value explicitly determined by $(\mathbf{h}, \mathbf{n})!$

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- This means any function that factors through *x* → (|*x_i* − *x_j*|_{*p*})_{*i*<*j*} is constant on *T_p*(𝔥, *n*), with value explicitly determined by (𝔥, *n*)!
- In particular, the product of the factors $|x_i x_j|_p$ has a nice form:

Key Fact 1:

Every $\mathbf{x} \in \mathcal{T}_p(\mathbf{fh}, \mathbf{n})$ satisfies $\prod_{i < j} |x_i - x_j|_p = p^{\binom{9}{2}} \prod_{\ell=0}^{3} p^{-\left[\sum_{\lambda \in \mathbf{fh}_{\ell}} \binom{\#\lambda}{2}\right]n_{\ell}} = p^{-29}$

The measure of $\mathcal{T}_p(\mathbf{fh}, \boldsymbol{n})$

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We attach a polynomial $M_{\Uparrow}(t) \in \mathbb{Z}[t]$ to \Uparrow using falling factorials:

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Key Fact 2:

The set $\mathcal{T}_{\rho}(\mathbf{fh}, \mathbf{n})$ is compact and open with Haar measure

$$M_{
m ft}(p) \cdot \prod_{\ell=0}^{3} p^{-{
m rank}({
m ft}_{\ell})n_{\ell}} = (p-1)^{6}(p-2)^{2} \cdot p^{-27}$$

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For each partition \pitchfork and $\beta \in \mathbb{C}$ it is convenient to define

$$E_{\pitchfork}(eta) := \mathsf{rank}(\pitchfork) + \sum_{\lambda \in \pitchfork} inom{\#\lambda}{2}eta,$$

for then if $Re(\beta)$ is sufficiently large we have...

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Key Fact 1 + Key Fact 2 \implies

$$\begin{split} \sum_{\boldsymbol{n}\in\mathbb{Z}_{>0}^4} \int_{\mathcal{T}_p(\boldsymbol{\pitchfork},\boldsymbol{n})} \prod_{i< j} |x_i - x_j|_p^\beta \, d\boldsymbol{x} &= \sum_{\boldsymbol{n}\in\mathbb{Z}_{>0}^4} p^{\binom{9}{2}\beta} \cdot M_{\boldsymbol{\pitchfork}}(p) \cdot \prod_{\ell=0}^3 p^{-E_{\boldsymbol{\pitchfork}_\ell}(\beta)n_\ell} \\ &= p^{\binom{9}{2}\beta} \cdot M_{\boldsymbol{\pitchfork}}(p) \cdot \prod_{\ell=0}^3 \frac{1}{p^{E_{\boldsymbol{\pitchfork}_\ell}(\beta)} - 1} \end{split}$$

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A tuple $\mathbf{h} = (\mathbf{h}_0, \dots, \mathbf{h}_L)$ of partitions of $\{1, 2, \dots, N\}$ is called a **splitting chain** of order N and length $L(\mathbf{h}) = L$ if

$$\{1, 2, \dots, N\} = h_0 > h_1 > \dots > h_L = \{1\}\{2\}\dots\{N\}$$

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an associated rational expression

$$J_{ightarrow t}(eta):= \mathit{M}_{ightarrow}(t)\cdot \prod_{\ell=0}^{\mathit{L}(ightarrow)-1}rac{1}{t^{\mathcal{E}_{ightarrow t}}(eta)-1}\in \mathbb{Q}(t,t^eta)$$

Theorem (W., 2020)

The p-adic Mehta integral converges for $\operatorname{Re}(\beta) > -2/N$ with value

$$\mathcal{Z}_{N}(\beta) = \left(\sum_{m \in \mathbb{Z}} \rho(p^{m}) p^{m(N + \binom{N}{2}\beta)}\right) \cdot \left(p^{\binom{N}{2}\beta} - p^{-N}\right) \cdot \sum_{\mathfrak{h} \in \mathcal{S}_{N}} J_{\mathfrak{h},p}(\beta)$$

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Note: The same strategy yields a more general formula for

$$\int_{\mathcal{K}^N} \rho(\|\boldsymbol{x}\|) \big(\max_{i < j} |x_i - x_j| \big)^{\boldsymbol{a}} \big(\min_{i < j} |x_i - x_j| \big)^{\boldsymbol{b}} \prod_{i < j} |x_i - x_j|^{\boldsymbol{s}_{ij}} \, d\boldsymbol{x}$$

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where K is an arbitrary nonarchimedean local field. This provides the canonical partition function for mixed-charge gases and joint moments of the diameter $\max_{i < j} |x_i - x_j|$ and minimal particle spacing $\min_{i < j} |x_i - x_j|$.

Examples: N = 2, 3, 4

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$$\sum_{\mathbf{h}\in\mathcal{S}_2}J_{\mathbf{h},t}(\beta)=\frac{t-1}{t^{1+\beta}-1}$$

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$$\sum_{\mathbf{h}\in\mathcal{S}_4} J_{\mathbf{h},t}(\beta) = \frac{(t-1)(t-2)(t-3)}{t^{3+6\beta}-1} + 4 \cdot \frac{(t-1)^2(t-2)}{(t^{3+6\beta}-1)(t^{2+3\beta}-1)} + 6 \cdot \frac{(t-1)^2(t-2)}{(t^{3+6\beta}-1)(t^{1+\beta}-1)}$$

$$egin{aligned} &+3\cdotrac{(t-1)^3}{(t^{3+6eta}-1)(t^{2+2eta}-1)}+6\cdotrac{(t-1)^3}{(t^{3+6eta}-1)(t^{2+2eta}-1)(t^{1+eta}-1)}\ &+12\cdotrac{(t-1)^3}{(t^{3+6eta}-1)(t^{2+3eta}-1)(t^{1+eta}-1)} \end{aligned}$$

A scary theorem and a quadratic recurrence

$$\#\mathcal{S}_N = \Omega\left(rac{(N!)^2}{(2\ln(2))^N\cdot N^{1+\ln(2)/3}}
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$$F_{N}(r,\beta) := \sum_{k=1}^{N-1} \frac{k}{N} \cdot \frac{\sinh\left(\frac{r}{2}\left[\left(N + \binom{N}{2}\beta\right)\left(1 - \frac{2k}{N}\right) + 1\right]\right)}{\sinh\left(\frac{r}{2}\left[\left(N + \binom{N}{2}\beta\right) - 1\right]\right)} \cdot F_{k}(r,\beta) \cdot F_{N-k}(r,\beta)$$

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- Note: If N and β are fixed, $r \mapsto F_N(r, \beta)$ is even and smooth

An efficient formula

Recall: The *p*-adic Mehta Integral with $N \ge 2$ and $\rho = \mathbf{1}_{[0,1]}$ has the form

$$\mathcal{Z}_{N}(eta) = \int_{\mathbb{Z}_{p}^{N}} \prod_{i < j} |x_{i} - x_{j}|_{p}^{eta} doldsymbol{x}$$

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Theorem (Sinclair and W., 2021)

The value of the integral can be computed efficiently via

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Grand canonical partition functions

• Suppose the gas also exchanges particles with the reservoir with "fugacity parameter" f

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- N is no longer constant, so we replace Z_N by a grand canonical partition function, defined for β ≥ 0 and f ∈ C by

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• Similarly, for log-Coulomb gases in $p\mathbb{Z}_p$ we define

$$\mathcal{Z}^{\circ}(\beta, f) := \sum_{N=0}^{\infty} \mathcal{Z}^{\circ}_{N}(\beta) \frac{f^{N}}{N!} \quad \text{where} \quad \mathcal{Z}^{\circ}_{N}(\beta) := \int_{(p\mathbb{Z}_{p})^{N}} \prod_{i < j} |x_{i} - x_{j}|_{p}^{\beta} d\boldsymbol{x}$$

Theorem (Sinclair, 2020)

The grand canonical partition function for log-Coulomb gas in \mathbb{Z}_p satisfies

$$\mathcal{Z}(\beta, f) = (\mathcal{Z}^{\circ}(\beta, f))^{p}$$

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The projective line $\mathbb{P}^1(\mathbb{Q}_p)$ is a compact metric space with...

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Definition (The projective *p*-adic Mehta Integral)

The canonical partition function for an N-particle log-Coulomb gas in $\mathbb{P}^1(\mathbb{Q}_p)$ is given by

$$\mathcal{Z}_{N}^{*}(\beta) := \int_{(\mathbb{P}^{1}(\mathbb{Q}_{p}))^{N}} \prod_{i < j} \delta(x_{i}, x_{j})^{\beta} d\mu^{N}$$

Abbreviated Theorem (W., 2021)

The integral $\mathcal{Z}_N^*(\beta)$ converges absolutely if and only if $\operatorname{Re}(\beta) > -2/N$. Like $\mathcal{Z}_N(\beta)$, it is a finite sum over splitting chains of order N. Each summand is a rational function of p and $p^{-\beta}$ closely resembling $J_{\oplus,p}(\beta)$.

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Note: This is a special case of a general formula for the integral

$$\int_{(\mathbb{P}^1(\mathcal{K}))^N} \prod_{i < j} \delta(x_i, x_j)^{s_{ij}} d\mu^N$$

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with K any nonarchimedean local field. It converges absolutely for precisely the same $s_{ij} \in \mathbb{C}$ as the integral $\int_{K^N} \rho(||\mathbf{x}||) \prod_{i < j} |x_i - x_j|^{s_{ij}} d\mathbf{x}$, and the set of such s_{ij} does not depend on K.

The projective analogue: An efficient formula

The functions F_N from before are also useful in the projective case:

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If $\operatorname{Re}(\beta) > -2/N$, the value of $\mathcal{Z}^*_N(\beta)$ can be computed efficiently via

$$\mathcal{Z}_{N}^{*}(\beta) = N! \sum_{k=0}^{N} \frac{\cosh\left(\frac{\log(p)}{2}\left(N + \binom{N}{2}\beta\right)\left(1 - \frac{2k}{N}\right)\right)}{\left(2\cosh\left(\frac{\log(p)}{2}\right)\right)^{N}} \cdot F_{k}(\log(p), \beta) \cdot F_{N-k}(\log(p), \beta)$$

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and $\mathcal{Z}_{N}^{*}(\beta)|_{p\mapsto p^{-1}} = \mathcal{Z}_{N}^{*}(\beta).$

The (p+1)th Power Law

There is also a grand canonical partition function for log-Coulomb gas in $\mathbb{P}^1(\mathbb{Q}_p)$:

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